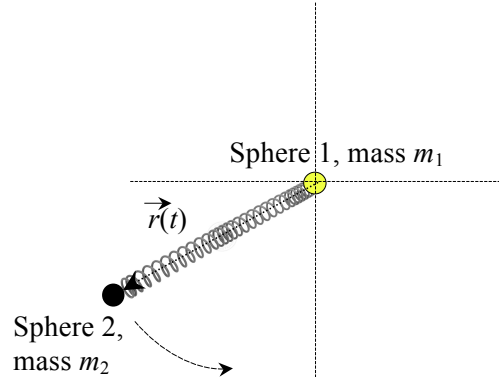


## P321b Midterm Practice Problem SOLUTIONS

### 1. Two spheres connected by a spring



Two spheres, of masses  $m_1$  and  $m_2$  respectively, are connected by a spring with spring constant  $k$  (and with zero length when unextended, so that the potential energy of the spring when stretched to length  $r$  is  $+\frac{1}{2}kr^2$ ). The spheres are orbiting around each other in a vacuum, so neither one of the spheres has a fixed position in space. Let us, however, consider a coordinate system which is non-rotating, but is always centred on Sphere 1 (i.e., rather than a coordinate system centred at the centre of mass), and thus we will only consider the position  $\vec{r}(t)$  of Sphere 2 *relative* to that of Sphere 1.

- (a) What is the Lagrangian for the full system, in terms of  $m_2$ ,  $\vec{r}(t)$ ,  $\dot{\vec{r}}(t)$ , and the spring constant  $k$ ?

$$L = \frac{1}{2} m_2 \dot{\vec{r}}^2 - \frac{1}{2} k r^2 .$$

- (b) Write down the resulting Lagrange's equations of motion, in terms of the polar coordinates  $r$  and  $\phi$  (in the plane of motion of the system).

In polar coordinates, we have  $L = \frac{1}{2} m_2 (\dot{r}^2 + r^2 \dot{\phi}^2) - \frac{1}{2} k r^2$ . Thus, we have  $d/dt(\partial L/\partial \dot{\phi}) = 0 \Rightarrow m_2 r^2 \dot{\phi} = \text{a constant angular momentum } \Lambda$ , and  $d/dt(\partial L/\partial \dot{r}) - (\partial L/\partial r) = 0 \Rightarrow m_2 \ddot{r} = m_2 r \dot{\phi}^2 - kr \Rightarrow \ddot{r} = -kr/m_2 + \Lambda^2/m_2^2 r^3$ .

- (c) What are the solutions to these equations of motion?

The differential equation above is not easy to solve. It is a complicated integral, and even though it doesn't look too horrible, it can't even be solved in Mathematica or Maple. So please don't worry about solving it – I certainly wouldn't ask you to solve a nasty differential equation like that on a real exam! But I will do it here anyway:

$\ddot{r} = d\dot{r}/dr \cdot dr/dt = \dot{r} \cdot d\dot{r}/dr = \alpha r + \beta/r^3$ , where  $\alpha = -k/m_2$  and  $\beta = \Lambda^2/m_2^2$ . Multiplying both sides by  $dr$  and integrating, we get that  $\frac{1}{2} \dot{r}^2 = \alpha r^2/2 - \beta/(2r^2) + C$ , where  $C$  is a constant.

Multiplying both sides by two, then square-rooting, and then multiplying both sides by  $dt$  gives us that  $t$  is equal to the integral of  $dr/\sqrt{(\alpha r^2 - \beta/r^2 + C)}$  plus another constant of integration  $D$ . Substituting  $s = r^2$  and  $ds = 2rdr$  gives us that  $t$  is equal to the integral of  $ds/(2\sqrt{(\alpha s^2 + Cs - \beta)}) + D$ . This integral can be found in big tables of integrals (like Gradshteyn & Ryzhik), and the result is that  $-1/(2\sqrt{\beta}) \arcsin((C - 2\beta s)/(C^2 - 4\alpha\beta)) = t + D$ . Substituting  $r$  back in and simplifying, one gets that  $r(t) = \sqrt{(\sqrt{m_2 C^2 - 4k\Lambda} \sin((2\Lambda t/m_2) + D) + C) / (\Lambda \sqrt{2})}$ . Ugh!

(d) What is the total energy of the system?

The total energy  $E = \dot{r}(\partial L/\partial \dot{r}) + \dot{\phi}(\partial L/\partial \dot{\phi}) - L \Rightarrow E = \frac{1}{2}m_2(\dot{r}^2 + r^2\dot{\phi}^2) + \frac{1}{2}kr^2$  (= T + U). Note that  $E$  can also be written in terms of the (constant) angular momentum as  $\frac{1}{2}m_2\dot{r}^2 + \Lambda^2/2m_2r^2 + \frac{1}{2}kr^2$ .

2. Two gravitating masses  $m_1$  and  $m_2$  are separated by a distance  $r_0$  and released from rest. Write down the Lagrangian and solve Lagrange's equations of motion to show that when the separation is  $r$  ( $< r_0$ ), the speeds are:  $v_1 = m_2 \sqrt{[2G/(m_1+m_2)](1/r - 1/r_0)}$  and  $v_2 = m_1 \sqrt{[2G/(m_1+m_2)](1/r - 1/r_0)}$ .

If  $x_1$  is the position of the first mass and  $x_2$  is the position of the second, the Lagrangian  $L = T - U = \frac{1}{2}m_1\dot{x}_1^2 + \frac{1}{2}m_2\dot{x}_2^2 + Gm_1m_2/(x_2 - x_1)$ . We also know that  $r = x_2 - x_1$  and, if we are working in the centre of mass frame, that  $m_1x_1 + m_2x_2 = 0$ . Thus  $x_1 = -m_2x_2/m_1$  and  $x_2 = r/(1 + m_2/m_1)$ , and the Lagrangian can be written purely in terms of  $r$ :

$L = m_1m_2\dot{r}^2/(m_1+m_2) + Gm_1m_2/r$ . Thus we have Lagrange's equation of motion for  $r$ :  $m_1m_2\ddot{r}/(m_1+m_2) = -Gm_1m_2/r^2 \Rightarrow \ddot{r} = -G(m_1+m_2)/r^2$ . One can obtain the solution to that equation of motion via  $\ddot{r} = \dot{r}d\dot{r}/dr = -G(m_1+m_2)/r^2$ , multiplying both sides by  $dr$ , and integrating, which gives you (defining  $v \equiv \dot{r}$ ) that  $v^2 = 2G(m_1+m_2)(1/r) + C$ . We know  $v^2 = 0$  when  $r = r_0$ , thus  $v^2 = 2G(m_1+m_2)(1/r - 1/r_0)$ , and thus  $v = \sqrt{2G(m_1+m_2)(1/r - 1/r_0)}$ . We also know that  $x_2 = m_1r/(m_1 + m_2)$ , thus  $v_2 = m_1v/(m_1 + m_2)$ , and thus  $v_2 = m_1 \sqrt{[2G/(m_1+m_2)](1/r - 1/r_0)}$ , and similarly, speed  $v_1 = m_2 \sqrt{[2G/(m_1+m_2)](1/r - 1/r_0)}$ .