PHYS506B Asignment #3 SOLUTIONS

1Consider the Dirac spinor field ψ. Let

15pts
$$
\Psi_L = P_L \Psi
$$
 $\overline{\Psi}_L = \overline{(\Psi_L)} \neq (\overline{\Psi})_L$
 $\Psi_R = P_R \Psi$ $\overline{\Psi}_R = \overline{(\Psi_R)} \neq (\overline{\Psi})_R$

Show that

a) $\psi \psi = \psi_{\textrm{\tiny L}} \psi_{\textrm{\tiny R}} + \psi_{\textrm{\tiny R}} \psi_{\textrm{\tiny L}}$ b) $\left(\overline{\psi}_\text{\tiny L} \psi_\text{\tiny R}\right)^\dagger = \overline{\psi}_\text{\tiny R} \psi_\text{\tiny L}$ c) $\overline{\psi}\gamma^{\mu}\psi = \overline{\psi}_{\scriptscriptstyle\perp}\gamma^{\mu}\psi_{\scriptscriptstyle\perp} + \overline{\psi}_{\scriptscriptstyle\mathrm{R}}\gamma^{\mu}\psi_{\scriptscriptstyle\mathrm{R}}$ ${\overline \Psi} \gamma^\mu \psi = {\overline \Psi}_{\shortparallel} \gamma^\mu \psi_{\shortparallel} + {\overline \Psi}_{\shortparallel} \gamma^\mu \psi$

05/04/2005

where

$$
P_{L} \equiv \frac{1}{2} (1 - \gamma^{5})
$$

\n
$$
P_{R} \equiv \frac{1}{2} (1 + \gamma^{5})
$$

2From the lecture notes (Running Coupling Constant), obtain the equations

30pts
\na)
$$
\left[-\frac{\partial}{\partial t} + \beta(\alpha)\frac{\partial}{\partial \alpha}\right]R\left(\frac{\rho^2}{\mu^2}, \alpha\right) = 0
$$

\nb) $\frac{\partial \alpha(Q)}{\partial t} = \beta(\alpha(Q))$ and $\frac{\partial \alpha(Q)}{\partial \alpha} = \frac{\beta(\alpha(Q))}{\beta(\alpha)}$
\nc) $\left[-\frac{\partial}{\partial t} + \beta(\alpha)\frac{\partial}{\partial \alpha}\right]R(1, \alpha(Q)) = 0$

Estimate the energy scale *Q* at which QED and QCD coupling constants meet. (Compare your 25pts result with the Planck mass scale given by *G*^{−1/2}). **3**

Hint: use the renormalization point μ = $M_{\rm Z}$ with

$$
\alpha_s(M_Z) = 0.118
$$
 $\alpha(M_Z) = 128^{-1}$ $M_Z = 91.187$ GeV

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PHYS506B Asignment #3

given 15/03/2005 due05/04/2005

Consider the definition of α (Q): $\;$ In $\frac{\varrho^2}{\mu^2}=\int_{\alpha(\mu)}^{\alpha(\varrho)}\frac{{\rm d} x}{\beta(x)}$ 2 $\ln \frac{Q^2}{2} = \int^{\alpha(Q)} \frac{dx}{x}$ *x* α μ^- Jα(μ = $\int_{\alpha(\mu)}^{\alpha(\mathcal{Q})} \frac{1}{\beta}$ **4**30pts

a) Show that to leading order in $β(x)$, that is with $β(x)=-bx^2$, we obtain the following result for $α(Q)$:

$$
\alpha(Q)^{-1} - \alpha(\mu)^{-1} = b \ln \frac{\varrho^2}{\mu^2}
$$

b) Using the next to leading order in $\beta(x)$, that is with $\beta(x)=-bx^2(1+b^2x)$, obtain an expression for $\alpha(Q)$.

540pts Consider the experimental results $\alpha_{\rm s}(M_{\rm Z})$ = 0.118 and $M_{\rm Z}$ = 91.187 GeV. Evolve this result for $\alpha_{\rm s}$ down to the scales *Q* ⁼ 35 GeV and *Q* ⁼ 2 GeV

a) using the leading order running of α_s ;

b) using the next to leading order for the running of α_{s} .

Hint: assume only one sharp mass threshold at the bottom quark mass, that is for each one of questions a) and b)

1) starting from the given $\alpha_{\rm s}(M_{\rm Z})$, using $\mu=M_{\rm Z}$, find $\alpha_{\rm s}$ at ${\sf Q}=35$ GeV and at ${\sf Q}=m_{\rm b}=4.3$ GeV;

2) starting from $\alpha_{\rm s}(m_{\rm b})$ found in 1), using $\mu=m_{\rm b}$, find $\alpha_{\rm s}$ at ${\sf Q}$ = 2 GeV.

Photocopy the figure on page 200 and plot your points on it.

N.B.: This treatment neglects a very small $\alpha_{_{\rm S}}$ shift at $m_{_{\rm D}}$ when going from ${\sf n}_{_{\sf f}}$ =5 to ${\sf n}_{_{\sf f}}$ =4.

Question 1

From
$$
P_L = \frac{1}{2}(1-\gamma^2)
$$
 $P_R = \frac{1}{2}(1+\gamma^2)$
\nand resulting that
\n $(\gamma^2)^2 = I$
\n $\gamma^4 \gamma^5 \gamma^4 = -(y^2)^1 = -\gamma^5$
\n $\left[\gamma^5 \gamma^{\mu}\right]_{+} = 0$
\nWe obtain
\n $P_L P_R = \frac{1}{\gamma}(1-\gamma^2)(1+\gamma^3) = \frac{1}{\gamma}(1-(\gamma^2)^2) = 0$
\n $P_R P_L = \frac{1}{\gamma}(1+\gamma^5)(1-\gamma^5) = \frac{1}{\gamma}(1-(\gamma^5)^2) = 0$
\n $P_R P_L = \frac{1}{\gamma}(1+\gamma^5)(1-\gamma^5) = \frac{1}{\gamma}(1-(\gamma^5)^3) = 0$
\n $P_R P_L = \frac{1}{2}(1+(\gamma^5)^+) = \frac{1}{2}(1+\gamma^5) = P_L$
\n $P_L P_L = \frac{1}{2}(1-(\gamma^5)^+) = \frac{1}{2}(1-\gamma^5) = P_L$
\n $P_L P_L = \frac{1}{2}(1-\gamma^5)\gamma^{\mu} = \frac{1}{2}(\gamma^{\mu} - \gamma^5)\gamma^{\mu}$
\n $= \frac{1}{2}(\gamma^{\mu} + \gamma^{\mu}\gamma^5) = \gamma^{\mu} \frac{1}{2}(1+\gamma^5) = \gamma^{\mu} P_R$

We Then Istain

 $42 = 24$ $48 = 24$ $\overline{\mathcal{H}} = (R \psi)^{\dagger} \gamma^{\circ} = \psi^{\dagger} R \gamma^{\circ} = \overline{\psi} R$ $\overline{\psi}_{R} = (P_{R}\psi)^{\dagger} \gamma^{\circ} = \psi^{\dagger} P_{R} \gamma^{\circ} = \overline{\psi} P_{L}$ $\psi = \psi_R + \psi_L$ $\overline{\psi} = \overline{\psi_R} + \overline{\psi_L}$

a) $\overline{Y} \Psi = (\overline{\Psi_1} + \overline{\Psi_2})(\Psi_1 + \Psi_3) = \Psi_1 \Psi_1 + \overline{\Psi_1} \Psi_3 + \Psi_5 \Psi_1 + \overline{\Psi_8} \Psi_8$ But $\vec{\psi}_L \psi_L = \vec{\psi} F_R F_L \psi = 0$ $\overline{\psi}_{R}\psi_{R} = \overline{\psi}_{R}R\Psi = 0$ $\overline{\psi}_{k}\psi_{R}=\overline{\psi}_{R}\psi_{k}\psi\neq0$ $\bar{\psi}_R \psi_L = \bar{\psi}_R R \psi \neq 0$ S $\overline{\psi}\psi = \overline{\psi} \psi_{\mathsf{k}} + \overline{\psi}_{\mathsf{k}} \psi_{\mathsf{L}}$ b) $(\overline{\psi_L}\psi_R)^{\dagger} = \psi_R^{\dagger}(\overline{\psi_L})^{\dagger} = \psi_R^{\dagger}(\psi_L^{\dagger}r^{\dagger})^{\dagger} = \psi_R^{\dagger}r^{\circ}\psi_L$ $=\bar{\psi}_{e}\psi_{L}$ $Q \Psi^{\mu\nu} = (\bar{H} + \bar{H}_{R})^{\nu} (H + H_{R})$ = $\overline{4}R\gamma^m4L + \overline{4}R\gamma^m4R + \overline{4}R\gamma^m4R + \overline{4}R\gamma^m4L$ $\omega J - \overline{\psi}_L \gamma^{\mu} \gamma^{\nu} \epsilon = \overline{\psi} \rho_{\mathcal{R}} \gamma^{\mu} \rho_{\mathcal{R}} \psi = \overline{\psi} \gamma^{\mu} \rho_{\mathcal{L}} \rho_{\mathcal{R}} \psi = 0$ $\bar{\psi}_{R}y^{\mu}\psi_{L} = \bar{\psi}R\psi - P_{L}\psi - \bar{\psi}y^{\mu}P_{R}P_{L}\psi = 0$ $\overline{\psi}_{L}y^{\mu}\psi_{L}$ = $\overline{\psi}R\psi^{\mu}R\psi^{\nu}$ = $\overline{\psi}y^{\mu}R\psi^{\mu}R\psi^{\nu}$ $\overline{\psi}_{R} \gamma^{\mu} \psi_{R} = \overline{\psi} R \gamma^{\mu} R \psi = \overline{\psi} \gamma^{\mu} R R \psi \neq 0$ 50 $\overline{4}y\mu\psi = \overline{4}y\lambda\mu\psi + \overline{4}a\gamma\mu\psi$

Question 2

a) We have $\mu^2 \frac{d}{d\mu^2} R\left(\frac{Q^2}{\mu^2}, \alpha(\mu)\right) = \left[\mu^2 \frac{d}{d\mu^2} + \mu^2 \frac{d\alpha(\mu)}{d\mu^2} \frac{d}{d\alpha(\mu)}\right] R\left(\frac{Q^2}{\mu^2}, \alpha(\mu)\right) = 0$ Lst τ = $\ln \varrho_1^2$ $\beta(\alpha(\mu)) = \mu^2 \frac{\partial \alpha(\mu)}{\partial \mu^2}$ Then $\frac{\partial z}{\partial u^2} = -\frac{u^2}{\partial z} \frac{\partial^2}{\partial v^2} = -\frac{1}{u^2}$ $\frac{\partial R}{\partial x^2} = \frac{\partial R}{\partial x} \frac{\partial R}{\partial y^2} = -\frac{1}{2} \frac{\partial R}{\partial x}$ This gives $[-\frac{1}{2\tau} + \beta(\alpha(\mu))] \frac{1}{2\alpha(\mu)} R(\frac{\alpha^{2}}{\mu^{2}}, \alpha(\mu)) = 0$ $\left[-\frac{\partial}{\partial \tau} + \beta(\alpha(\omega)) \frac{\partial}{\partial \alpha(\omega)} \right] R(e^{\tau}, \alpha(\omega)) = 0$

b) Conoiden The definition of
$$
\alpha(\alpha)
$$

 $\tau(\alpha) = \int_{\alpha(\mu)}^{\alpha(\alpha)} \frac{dx}{\beta(\alpha)}$

We Then have

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\n
$$
\frac{\partial Z(G)}{\partial T} = 1 = \frac{1}{\beta(\alpha(G))} \frac{\partial \alpha(G)}{\partial C} - \frac{1}{\beta(\alpha(M))} \frac{\partial \alpha(M)}{\partial C}
$$
\nBut $\frac{\partial \alpha(M)}{\partial C} = 0$. Apply
\n
$$
\frac{\partial \alpha(G)}{\partial C} = \beta(\alpha(G))
$$

 \mathbf{r}_1 , \mathbf{r}_2 , \mathbf{r}_3 , \mathbf{r}_4 , \mathbf{r}_5

Also, we have

$$
\frac{\partial \tau(c)}{\partial \alpha(\mu)} = 0 = \frac{1}{\beta(\alpha(c))} \frac{\partial \alpha(\mu)}{\partial \alpha(\mu)} - \frac{1}{\beta(\alpha(\mu))} \frac{\partial \alpha(\mu)}{\partial \alpha(\mu)}
$$
\n
$$
\frac{\partial \alpha(\mu)}{\partial \alpha(\mu)} = 1 - \frac{\beta(\alpha(c))}{\beta(\alpha(\mu))}
$$
\n
$$
\frac{\partial \alpha(\mu)}{\partial \alpha(\mu)} = \frac{\beta(\alpha(c))}{\beta(\alpha(\mu))}
$$

c) With $R = R(1, \alpha(\varrho))$, we have $\left[-\frac{\partial}{\partial t} + \beta(\alpha(\mu))\frac{\partial}{\partial\alpha(\mu)}\right]R$ $= -\frac{\partial R}{\partial t} + \beta \left(\alpha(\mu) \right) \frac{\partial R}{\partial (\alpha|\mu|)}$ $= -\frac{\partial R}{\partial \alpha(Q)} \frac{\partial \alpha(Q)}{\partial Z} + \beta \left[\alpha(M) \right] \frac{\partial R}{\partial \alpha(Q)} \frac{\partial \alpha(Q)}{\partial \alpha(Q)}$ = $\frac{\partial R}{\partial \alpha(\alpha)}$ $\left[-\frac{\partial \alpha(\alpha)}{\partial \tau} + \beta(\alpha(\mu)) \frac{\partial \alpha(\alpha)}{\partial \alpha(\mu)} \right]$ = $\frac{\partial R}{\partial x(\alpha)}$ $\left[-\beta(x(\alpha)) + \beta(x(\mu)) \cdot \frac{\beta(x(\alpha))}{\beta(x(\mu))} \right]$ $= 0$

Question 38
\nQuestion 38
\n1 a
$$
tan\theta
$$
 or $cos\theta$, we should.
\n1 a $tan\theta$ or $cos\theta$ or

Also ses Holzen-Martin p348.

Question 4

 $\int_{0}^{1} \int_{0}^{x} f(x) dx = -bx^{2}$ $\alpha(G)$ $T = \int_{\alpha}^{\alpha(\epsilon)} \frac{dx}{\beta(\epsilon)}$ Then $\alpha(G)$ $-\frac{dx}{b}$ $\overline{\alpha}$ $\frac{1}{b}$ $\frac{1}{\alpha(\omega)}$ $\frac{1}{\alpha_{s}}$ This spelds $\alpha(\alpha) - \alpha_{0}^{-1} = b \ln(\frac{a^{2}}{\mu^{2}})$ OR $\alpha(\alpha) = \frac{\alpha_0}{1 + \alpha_0 b \ln(\frac{Q_2}{M^2})}$ This can also be unter $\alpha(\alpha)$ - $\alpha'(\mu)$ = $b \ln (\alpha')$ a high is invariant under a $\leftrightarrow \mu$ Though The innung of X depends on The order of which B(x) is expected, it does not depend on the renormalization scale by definition We also veryly That $\alpha(Q-\mu) = \alpha(\mu)$

Indeed, we using That $\frac{\frac{1}{\alpha(n)} + b\ln(\frac{Q^{2}}{n^{2}})}{+ b\ln(\frac{Q^{2}}{n^{2}})}$ $\frac{\alpha(\mathbb{Q}_1)}{\alpha(\mathbb{Q}_2)}$ Let $\mu = \mathbb{Q}_3$; Then $\frac{1}{\alpha(G_3)} = \frac{1}{\alpha(G_2)} + b \ln \frac{G_3^2}{Q_2^2}$ $=\frac{1}{\alpha(G_1)}+bln \frac{Q_3^2}{Q_1^2}$ Margare $\frac{\alpha(Q_i)}{\alpha(Q_1)}(\mu = Q_3) = \frac{\frac{1}{\alpha(Q_1)} + b \ln \frac{Q_2^2}{Q_3^2} \cdot \frac{Q_3^2}{Q_2^2}}{\frac{1}{\alpha(Q_1)} + b \ln \frac{Q_1^2}{Q_3^2} \cdot \frac{Q_2^2}{Q_1^2}}$ $\frac{\alpha(G)}{\alpha(G)}$ Me see that indeed the unning of α does not depend on u.

 $= -bx^{2} - bb'x^{3} = -b'x^{2}(1+b'x)$ Let $\beta(x)$ A
 $T = \int_{\alpha}^{\alpha(\mathbb{G})} \frac{dx}{\beta(x)} = \int_{\alpha}^{\alpha(\mathbb{G})} \frac{dx}{\beta(x)}$ 1km $\frac{dx}{x^{2}(ax+b)} = -\frac{1}{b^{2}} + \frac{a}{b^{2}} ln$ \overline{r} $\left(\frac{ax+b}{x}\right)$ Therefore $- bT = \left[-\frac{1}{x} + b' \ln \left(\frac{1+b'x}{x} \right) \right]_{\alpha}^{\alpha(\alpha)}$ = $-\frac{1}{\alpha(G)} + \frac{1}{\alpha_{0}} + b'h(1+b'\alpha(G)) - b'h(1+b'\alpha_{0})$ atrich can be uniten as $a^{-1}(a) - a^{-1}(\mu) + b^{1}ln\left(\frac{\alpha(a)}{\alpha(a)}\right) - b^{1}ln\left(\frac{\alpha(\mu)}{\alpha + b^{1}(\alpha(\mu))}\right) = bln \frac{a^{2}}{\mu^{2}}$ note That This is invariant under $Q \Leftrightarrow \mu$ We also very that indeed $\alpha(\alpha-\mu) = \alpha(\mu)$! More elegantly: $\alpha^{-1}(a) - \alpha'(n) + b' ln \left[\frac{\alpha(a)}{\alpha(n)} \frac{1+b' \alpha(n)}{1+b' \alpha(a)} \right] = b ln \frac{a}{n^2}$

Question 5
$$
\frac{1}{L_1C_1}
$$
 $\frac{1}{\alpha} - (\alpha) = \frac{1}{\alpha} - (\alpha) + b/\ln \frac{\alpha}{\mu}$ = c
\n*N.L.O.* $\frac{1}{\alpha} - (\alpha) = \frac{1}{\alpha} - (\mu) + b/\ln \frac{\alpha}{\mu}$ $\frac{1}{\alpha}(\frac{1}{\alpha}) = \frac{1}{\alpha} - \frac{1}{\alpha}(\frac{1}{\alpha}) - \frac{1}{\alpha} - \frac{1}{\alpha}(\frac{1}{\alpha})$
\nb $= \frac{1}{\frac{1}{\alpha} - (\frac{1}{\alpha}) - \frac{1}{\alpha} - \frac{1}{\alpha}$

 $\frac{1}{\sqrt{2}}$

L

 \pm 1₀ limits of the PDG $\frac{3}{50}$
average The figure central values and the values of $\alpha_s(Q)$ at the
values of Q where average. The figure they are measured. The lines show the decrease in α_s with clearly shows the Summary of the increasing Q.

PHYS506B, spring 2003

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Part I: Introduction to QCD