## PHYS506B Asignment #3 SOLUTIONS

given 15/03/2005 due 05/04/2005

1 Consider the Dirac spinor field  $\psi$ . Let

15pts

$$\begin{array}{ll} \psi_{\text{\tiny L}} = P_{\text{\tiny L}} \psi & \overline{\psi}_{\text{\tiny L}} \equiv \overline{\left( \psi_{\text{\tiny L}} \right)} \neq \left( \overline{\psi} \right)_{\text{\tiny L}} & P_{\text{\tiny L}} \equiv \frac{1}{2} \left( 1 - \gamma^5 \right) \\ \psi_{\text{\tiny R}} = P_{\text{\tiny R}} \psi & \overline{\psi}_{\text{\tiny R}} \equiv \overline{\left( \psi_{\text{\tiny R}} \right)} \neq \left( \overline{\psi} \right)_{\text{\tiny R}} & \text{where} & P_{\text{\tiny R}} \equiv \frac{1}{2} \left( 1 + \gamma^5 \right) \end{array}$$

Show that

a) 
$$\overline{\psi}\psi = \overline{\psi}_L\psi_R + \overline{\psi}_R\psi_L$$

b) 
$$(\overline{\psi}_L \psi_R)^{\dagger} = \overline{\psi}_R \psi_L$$

c) 
$$\overline{\psi}\gamma^{\mu}\psi = \overline{\psi}_{L}\gamma^{\mu}\psi_{L} + \overline{\psi}_{R}\gamma^{\mu}\psi_{R}$$

2 From the lecture notes (Running Coupling Constant), obtain the equations

30pts

a) 
$$\left[ -\frac{\partial}{\partial t} + \beta(\alpha) \frac{\partial}{\partial \alpha} \right] R\left(\frac{Q^2}{\mu^2}, \alpha\right) = 0$$
 b)  $\frac{\partial \alpha(Q)}{\partial t} = \beta(\alpha(Q))$  and  $\frac{\partial \alpha(Q)}{\partial \alpha} = \frac{\beta(\alpha(Q))}{\beta(\alpha)}$ 

c)  $\left[ -\frac{\partial}{\partial t} + \beta(\alpha) \frac{\partial}{\partial \alpha} \right] R(1, \alpha(Q)) = 0$ 

**3** Estimate the energy scale Q at which QED and QCD coupling constants meet. (Compare your result with the Planck mass scale given by  $G^{-1/2}$ ).

Hint: use the renormalization point  $\mu = M_7$  with

$$\alpha_{s}(M_{z}) = 0.118$$
  $\alpha(M_{z}) = 128^{-1}$   $M_{z} = 91.187 \text{ GeV}$ 

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- 4 Consider the definition of  $\alpha(Q)$ :  $\ln \frac{Q^2}{\mu^2} = \int_{\alpha(\mu)}^{\alpha(Q)} \frac{dx}{\beta(x)}$ 
  - a) Show that to leading order in  $\beta(x)$ , that is with  $\beta(x) = -bx^2$ , we obtain the following result for  $\alpha(Q)$ :

$$\alpha(Q)^{-1} - \alpha(\mu)^{-1} = b \ln \frac{Q^2}{\mu^2}$$

- b) Using the next to leading order in  $\beta(x)$ , that is with  $\beta(x) = -bx^2(1+b'x)$ , obtain an expression for  $\alpha(Q)$ .
- **5** Consider the experimental results  $\alpha_s(M_Z) = 0.118$  and  $M_Z = 91.187$  GeV. Evolve this result for  $\alpha_s$  40pts down to the scales Q = 35 GeV and Q = 2 GeV
  - a) using the leading order running of  $\alpha_s$ ;
  - b) using the next to leading order for the running of  $\alpha_{\rm s}$ .

Hint: assume only one sharp mass threshold at the bottom quark mass, that is for each one of questions a) and b)

- 1) starting from the given  $\alpha_s(M_Z)$ , using  $\mu=M_Z$ , find  $\alpha_s$  at Q=35 GeV and at  $Q=m_b=4.3$  GeV;
- 2) starting from  $\alpha_s(m_b)$  found in 1), using  $\mu=m_b$ , find  $\alpha_s$  at Q=2 GeV.

Photocopy the figure on page 200 and plot your points on it.

N.B.: This treatment neglects a very small  $\alpha_s$  shift at  $m_h$  when going from  $n_f$ =5 to  $n_f$ =4.

From 
$$P_L = \frac{1}{2}(1-r^s)$$
  $P_R = \frac{1}{2}(1+r^s)$  and recalling that

 $(r^s)^2 = I$ 
 $y^ay^sy^a = -(y^s)^{\frac{1}{2}} = -y^s$ 
 $\begin{bmatrix} y^s, y^{\mu} \end{bmatrix}_+ = 0$ 

We obtain

 $P_L P_R = \frac{1}{4}(1-r^s)(1+r^s) = \frac{1}{4}(1-(r^s)^2) = 0$ 
 $P_R P_L = \frac{1}{4}(1+r^s)(1-r^s) = \frac{1}{4}(1-(r^s)^2) = 0$ 
 $P_R P_L = \frac{1}{4}(1+r^s)(1-r^s) = \frac{1}{4}(1-r^s)^2 = P_R$ 
 $P_L^+ = \frac{1}{2}(1-(r^s)^+) = \frac{1}{2}(1-r^s) = P_L$ 
 $P_L^+ = \frac{1}{2}(1-(r^s)^+) = \frac{1}{2}(1-r^s) = P_L$ 
 $P_L^+ = \frac{1}{2}(1-r^s)r^m = \frac{1}{2}(r^m-r^s)r^m$ 
 $= \frac{1}{2}(r^m+r^mr^s) = r^m\frac{1}{2}(1+r^s) = r^mP_R$ 

We Then obtain

 $Y_L = P_L Y$ 
 $Y_R = (P_R Y_L)^+r^0 = Y_R Y_R^0 = \overline{Y}_R^0$ 
 $\overline{Y}_R = (P_R Y_L)^+r^0 = Y_R^0$ 
 $\overline{Y}_R = (P_R Y_L)^+r^0 = Y_R^0$ 
 $\overline{Y}_R = (P_R Y_L)^+r^0 = Y_R^0$ 
 $\overline{Y}_R = (P_R Y_L)^+r^0 = Y_R^0$ 

a) 
$$\Psi\Psi = (\Psi_L + \Psi_R)(\Psi_L + \Psi_R) = \Psi_L\Psi_L + \Psi_L\Psi_R + \Psi_R\Psi_L + \Psi_R\Psi_R + \Psi_R\Psi$$

b) 
$$(\bar{\Psi}_L \Psi_R)^{\dagger} = \Psi_R^{\dagger} (\bar{\Psi}_L)^{\dagger} = \Psi_R^{\dagger} (\Psi_L^{\dagger} Y^{\circ})^{\dagger} = \Psi_R^{\dagger} Y^{\circ} \Psi_L$$

$$= \bar{\Psi}_R \Psi_L$$

Question 2

a) We home

$$\mu^{2} \frac{1}{d_{1}^{2}} R\left(\frac{Q^{2}}{M^{2}}, \alpha(M)\right) = \left[M^{2} \frac{1}{\partial M^{2}} + M^{2} \frac{1}{\partial \alpha(M)} \frac{1}{\partial M}\right] R\left(\frac{Q^{2}}{M^{2}}, \alpha(M)\right) = 0$$

Let  $t = \ln Q^{2}$ 

$$\beta(\alpha(M)) = \mu^{2} \frac{1}{\partial \alpha(M)}$$

$$\frac{1}{\partial M^{2}} = -\frac{M^{2}}{Q^{2}} \frac{Q^{2}}{M^{2}} = -\frac{1}{M^{2}}$$

$$\frac{1}{\partial M^{2}} = \frac{1}{\partial M^{2}} \frac{1}{\partial M^{2}} = -\frac{1}{M^{2}} \frac{1}{\partial T}$$
Thus gives
$$\left[-\frac{1}{\partial T} + \beta(\alpha(M)) \frac{1}{\partial \alpha(M)}\right] R\left(\frac{Q^{2}}{M^{2}}, \alpha(M)\right) = 0$$

$$\left[-\frac{1}{\partial T} + \beta(\alpha(M)) \frac{1}{\partial \alpha(M)}\right] R\left(e^{T}, \alpha(M)\right) = 0$$

b) Consider The definition of 
$$\alpha(Q)$$

$$T(Q) = \int_{\alpha(W)} \frac{dx}{\beta(x)}$$

$$\frac{\partial \mathcal{I}(G)}{\partial \mathcal{I}} = 1 = \frac{1}{\beta(\alpha(G))} \frac{\partial \alpha(G)}{\partial \mathcal{I}} - \frac{1}{\beta(\alpha(M))} \frac{\partial \alpha(M)}{\partial \mathcal{I}}$$
But  $\frac{\partial \alpha(M)}{\partial \mathcal{I}} = 0$ . Prespective 
$$\frac{\partial \alpha(G)}{\partial \mathcal{I}} = \beta(\alpha(G))$$

$$\frac{\partial \mathcal{L}(G)}{\partial \alpha(\mu)} = 0 = \frac{1}{\beta(\alpha(G))} \frac{\partial \alpha(G)}{\partial \alpha(\mu)} - \frac{1}{\beta(\alpha(\mu))} \frac{\partial \alpha(\mu)}{\partial \alpha(\mu)}$$

c) With 
$$R = R(1, \alpha(Q))$$
, we have
$$\begin{bmatrix} -\frac{1}{2} + \beta(\alpha(M)) \frac{1}{2\alpha(M)} \end{bmatrix} R$$

$$= -\frac{1}{2}R + \beta(\alpha(M)) \frac{1}{2}R \frac{1}{2}(\alpha(M))$$

$$= -\frac{1}{2}R + \beta(\alpha(M)) \frac{1}{2}R \frac{1}{2}(\alpha(M))$$

$$= -\frac{1}{2}R \frac{1}{2}\alpha(Q) \frac{1}{2}R + \beta(\alpha(M)) \frac{1}{2}R \frac{1}{2}\alpha(Q)$$

$$= \frac{1}{2}R \frac{1}{2}R \frac{1}{2}R - \beta(\alpha(Q)) + \beta(\alpha(M)) \frac{1}{2}R \frac{1}{2}$$

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QED and QCD.
  to do a proper jeb, we should.
      10 evalue of from Q=MZ to Q= Mtop using Ng=5 avalue of from Q=MZ to Q= Mtop using Ng=5
      2° Let \alpha(C_5) = \alpha^2(C_5) using ul = e
  But to obtain an estimate, we proceed with step 2° using n/26 starting from \alpha(m_2) and \alpha_s(m_2).
   at Lo us hore
  Q(Q) = < < (Q) 40 1 + b ln Q2 = 1 x (m) + bs Q2 x 2
         \ln \frac{G^2}{M^2} = \frac{1}{b_s - b} \left[ \frac{1}{x(m)} - \frac{1}{x_s(m)} \right]
        b = -\frac{1}{3\pi} \left| \sum_{i=1}^{n_b} e_i^2 + \sum_{i=1}^{n_b} 3 e_i^2 \right| where
            b = -\frac{1}{3\pi} \left[ 3 + 3^2 \left( \frac{2}{3} \right)^2 + \left( \frac{1}{3} \right)^2 \right] = -\frac{8}{3\pi}
           b_5 = \frac{1}{12\pi} (33 - 2n_0) = \frac{7}{4\pi}
                  Starting from
                       ~ (Mz) = 0,00781 = 128-1
                        MZ = 91,187 GeV.
                   In Q2 = 85,02 -0 Q2 = 85,0 36
Wo obstern
                   Q ~ 2.9 X10 -0 Q ~ 2.6 X10 GeV
                 G-1/2 = 1,22 X10 GOV
   har sterr
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Question 3

Alas see Holzen-Mantin p348.

Let u= Qz; Then  $\frac{1}{\alpha(G_3)} = \frac{1}{\alpha(G_2)} + b \ln \frac{G_3^2}{Q_2^2}$  $= \frac{1}{\alpha(\alpha_1)} + b \ln \frac{\alpha_3^2}{\alpha_1^2}$ Marefore  $\frac{\langle (Q_1) (\mu = Q_3) | - \frac{1}{\langle (Q_2) | + b \ln \frac{Q_2^2}{Q_3^2} \cdot \frac{Q_3^2}{Q_2^2}}{|Q_2|^2}$   $\frac{1}{\langle (Q_1) | + b \ln \frac{Q_1^2}{Q_3^2} \cdot \frac{Q_2^2}{Q_2^2}}{|Q_2|^2}$ We see that indeed the running of & does not depend on u.

Let 
$$\beta(x) = -bx^2 - bb x^3 = -bx^2 (1+b'x)$$

Then
$$\alpha(G)$$

$$T = \int_{\infty}^{\infty} \frac{dx}{\beta(x)} = \int_{\infty}^{\infty} -bx^2 (1+b'x)$$
but
$$\int_{\infty}^{\infty} \frac{dx}{\beta(x)} = \int_{\infty}^{\infty} -bx^2 (1+b'x)$$
but
$$\int_{\infty}^{\infty} \frac{dx}{\beta(x)} = \int_{\infty}^{\infty} -bx^2 (1+b'x)$$

$$\int_{\infty}^{\infty} \frac{dx}{\beta(x)} = \int_{\infty}^{\infty} \frac{dx}{\beta(x)} + \int_{\infty}^{\infty} \frac{dx}{\beta(x)} = \int_{\infty}^{\infty} \frac{dx}{\beta(x)}$$
Therefore
$$\int_{\infty}^{\infty} -bx^2 (1+b'x)$$

$$\int_{\infty}^{\infty} \frac{dx}{\beta(x)} = \int_{\infty}^{\infty} \frac{dx}{\beta(x)} + \int_{\infty}^{\infty} \frac{dx}{\beta(x)} = \int_{\infty}^{\infty} \frac{dx}{\beta(x)}$$

$$\int_{\infty}^{\infty} \frac{dx}{\beta(x)} = \int_{\infty}^{\infty} \frac{dx}{\beta(x)} + \int_{\infty}^{\infty} \frac{d$$

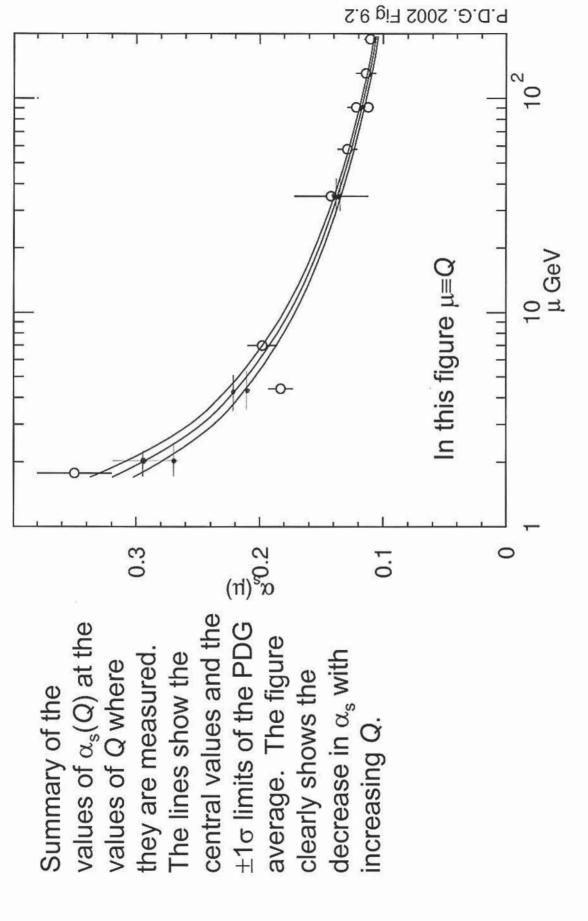
Question 5

L.O. 
$$\propto (G) - \alpha'(N) - b \ln \frac{G^2}{M^2} = 6$$

N.L.O.  $\propto (G) - \alpha'(N) + b \ln \frac{G^2}{M^2} = 6$ 

N.L.O.  $\propto (G) - \alpha'(N) + b \ln \frac{G^2}{M^2} = 6$ 
 $b = \frac{1}{12\pi} (33 - 2n)$ 
 $b' = \frac{153 - 19n!}{2\pi\pi} (33 - 2n)$ 
 $b' = \frac{23}{12\pi} = 0.600$ 
 $\frac{25}{12\pi} = 0.663$ 
 $\frac{29}{12\pi} = 0.401$ 
 $\frac{29}{23\pi} = 0.401$ 
 $\frac$ 

## Running Coupling Constant



100

Part I: Introduction to QCD

PHYS506B, spring 2003